

Asymptotic approximations for American options

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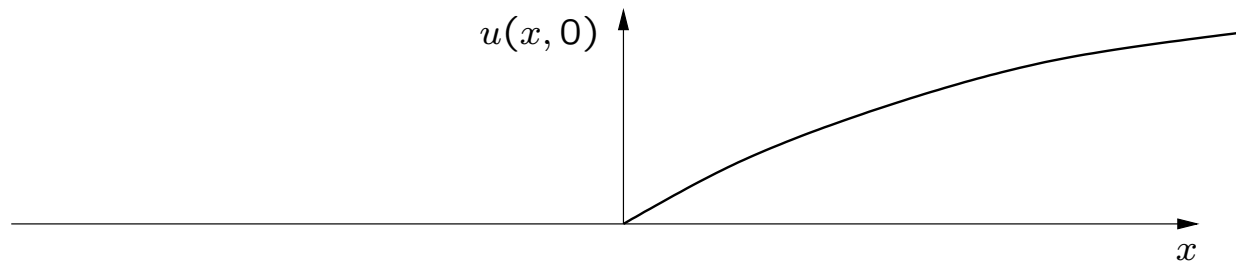
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Short-time asymptotics of the heat equation

Consider

$$u_t = \frac{1}{2}u_{xx}, \quad t > 0,$$

with initial data vanishing for $x < 0$:



Initial behaviour as $x \rightarrow -\infty$, $t = O(1)$? (Or as $t \downarrow 0$, x fixed.)

Various ways:

$$\begin{aligned}u(x, t) &= \frac{1}{\sqrt{2\pi t}} \int_0^\infty u_0(s) e^{-(x-s)^2/2t} ds \\&= \frac{e^{-x^2/2t}}{\sqrt{2\pi t}} \int_0^\infty u_0(s) e^{xs/t - s^2/2t} ds \quad (*) \\&= \frac{e^{-x^2/2t}}{\sqrt{2\pi t}} \left(\frac{-x}{t} \right) \int_0^\infty u_0(-t\xi/x) e^{-\xi - t\xi^2/2x^2} d\xi \\&\sim \frac{e^{-x^2/2t}}{\sqrt{2\pi t}} \times F(x/t) \quad \text{as } x \rightarrow -\infty;\end{aligned}$$

or do Laplace on (*) above,

$$\frac{e^{-x^2/2t}}{\sqrt{2\pi t}} \int_0^\infty u_0(s) e^{xs/t - s^2/2t} ds;$$

as $x/t \rightarrow -\infty$;

or expand $u_0(s) = \sum c_n s^n$ and get the answer as a sum of similarity solutions;

(both these conclude that the behaviour of u_0 at the origin is paramount)

or

$$u = \frac{e^{-x^2/2t}}{\sqrt{2\pi t}} v(x, t)$$

gives

$$tv_t + xv_x = \frac{1}{2}v_{xx}$$

so put RHS = 0 and say Euler; or put $x = X/\epsilon$
(or $t = \epsilon^2 T$) and use the WKB ansatz

$$u \sim Ae^{v/\epsilon^2}$$

to get the same result via $v_T = \frac{1}{2}v_X^2$ etc.

The Stefan problem for small latent heat

Melting of a solid with small latent heat ϵ :

$$u_t = \frac{1}{2}u_{xx}, \quad 0 < x < s(t), \quad u(0, t) = 1,$$

with free boundary conditions

$$u(s(t), t) = 0, \quad u_x(s(t), t) = -\epsilon \dot{s}.$$

There is a similarity solution $u = U(x/\sqrt{t})$,
 $s = \alpha\sqrt{t}$ from which, as $\epsilon \rightarrow 0$, the relevant
timescale is

$$t = \delta T, \quad \delta = 1/|\log \epsilon|.$$



Then there is a 3-layer structure:

Boundary layer near $x = 0$, $x = \delta^{\frac{1}{2}}X$, $t = \delta T$, giving the usual error function solution.

'Outer region' $x = O(1)$, $t = \delta T$, WKB solution (as above) of

$$\frac{1}{\delta}u_T = u_{xx}$$

Solution is $u \sim A \exp(v/\delta)$ with $v = -x^2/2T$,
 $A = (1/\sqrt{2\pi T})(T/x)$.

Inner layer $x = s(t) + \delta\xi$, $u = \epsilon U$, with a travelling wave solution of the heat equation satisfying the free boundary conditions, $U = \frac{1}{2}(1 - \exp(-2\xi\dot{s}))$.

These all match and the scale δ follows from matching the outer region to the inner layer. Generalises to more than one dimension and the free boundary is close to the isotherm

$u = \epsilon$ of the corresponding pure heat conduction problem. Can also be done via an integral equation (Grinberg/Chekhmareva) but doesn't work in 2 or more dimensions.

(Addison, SDH, King, QAM 2005.)

American options in the Black-Scholes model

The BS model is the standard description of normal (?!) financial markets.

- Asset prices follow diffusions (SDEs driven by Wiener processes).

- Options are contracts paying a given function $P(S_T)$, the *payoff*, of the asset price S_T on a final date $t = T$.
- Options are valued as expectations; by Feynman-Kac, option prices satisfy a backward parabolic equation in S, t , with final data $P(S)$: the

BS PDE

$$\frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - q)S \frac{\partial V}{\partial S} - rV = 0.$$

A simple scaling and time-reversal

$$t' = \sigma^2(T - t)$$

(so t' is dimensionless) turns

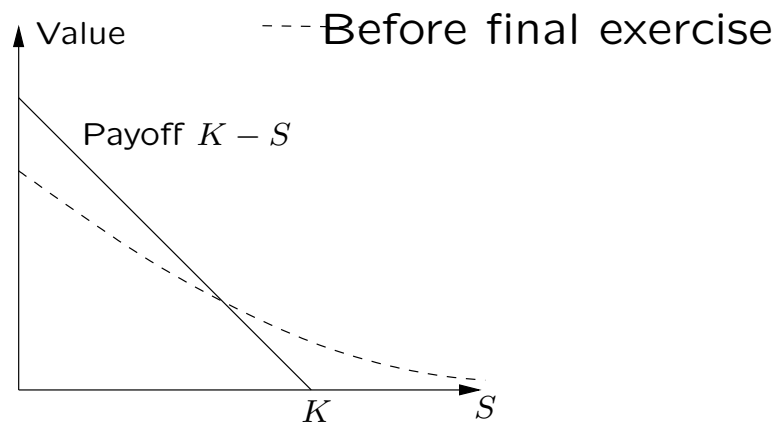
$$\frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - q)S \frac{\partial V}{\partial S} - rV = 0.$$

into

$$\frac{\partial V}{\partial t'} = \frac{1}{2}S^2 \frac{\partial^2 V}{\partial S^2} + (\rho - \gamma)S \frac{\partial V}{\partial S} - \rho V, \quad \rho = \frac{r}{\sigma^2}, \quad \gamma = \frac{q}{\sigma^2},$$

with the payoff as *initial* data.

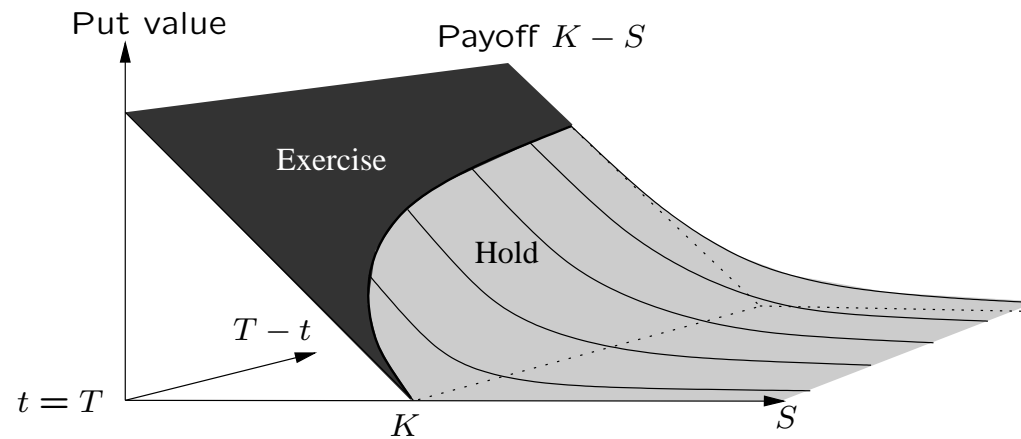
$$\frac{\partial V}{\partial t'} = \underbrace{\frac{1}{2} S^2 \frac{\partial^2 V}{\partial S^2}}_{\text{convexity increases } V} + \underbrace{(\rho - \gamma) S \frac{\partial V}{\partial S} - \rho V}_{\text{both decrease } V}$$



An American option can be exercised at any time (not just at the final date).

Hence **option value** \geq **payoff**.

The American option is like a continuous series of obstacle-type problems (a parabolic variational inequality).



Optimality \longrightarrow 'smooth pasting' free boundary conditions: V and $\partial V/\partial S$ are continuous at the interface $S = S^*(t)$:

$$V = K - S, \quad \frac{\partial V}{\partial S} = -1, \quad S = S^*(t).$$

Discrete dividend payments

When dividends are paid the asset price falls (in calendar time t):

$$S_{\text{before}} = S_{\text{after}} + \text{dividend}$$

The model above has dividends paid continuously at rate q , asset price process

$$\frac{dS_t}{S_t} = (r - q) dt + \sigma dW_t$$

The corresponding scaled and forwardised BS
PDE is

$$\frac{\partial V}{\partial t'} = \frac{1}{2}S^2 \frac{\partial^2 V}{\partial S^2} + (\rho - \gamma)S \frac{\partial V}{\partial S} - \rho V, \quad \rho = \frac{r}{\sigma^2}, \quad \gamma = \frac{q}{\sigma^2}.$$

For *discrete* dividends, paying $qS_{t_n^-} \delta t$ at (equal) time intervals t_n separated by δt ,

$$S_{t_n^+} = (1 - q \delta t) S_{t_n^-},$$

or in scaled time $T - t = \sigma^2 t'$,

$$S_{t'_n-} = (1 - \gamma \epsilon^2) S_{t'_n+}, \quad \boxed{\epsilon^2 = \sigma^2 \delta t.}$$

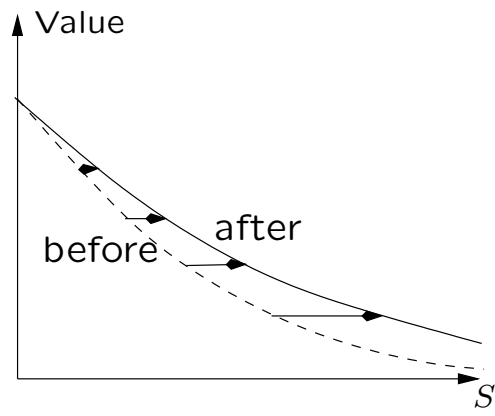
Between these dates, zero-dividend forwardised BS PDE holds:

$$\frac{\partial V}{\partial t'} = \frac{1}{2} S^2 \frac{\partial^2 V}{\partial S^2} + \rho S \frac{\partial V}{\partial S} - \rho V.$$

At dividend dates, **option value is continuous** for each realisation of S_t , so $V(S_{t'_n+}, t'_n+) = V(S_{t'_n-}, t'_n-)$ which is

$$V(S, t'_n+) = V((1 - \gamma\epsilon^2)S, t'_n-)$$

for all $0 < S < \infty$. That is, the option values are **shifted to the right** across a dividend date (in backwards time).



Discrete PDE + jump cond's to cont's PDE

Multiple scale ansatz $V(S, t', \tau)$ where

$$t' = t'_n + \epsilon^2 \tau$$

so discrete problem is

$$\frac{\partial V}{\partial t'} + \frac{1}{\epsilon^2} \frac{\partial V}{\partial \tau} = \frac{1}{2} S^2 \frac{\partial^2 V}{\partial S^2} + \rho S \frac{\partial V}{\partial S} - \rho V, \quad 0 < \tau < 1$$

with...

$$V(S, t', 1^+) = V((1 - \gamma\epsilon^2)S, t', 1^-)$$

and **periodic in** τ to eliminate secular terms,
so

$$V(S, t', 1^+) = V(S, t', 0^+).$$

Expand

$$V \sim V_0 + \epsilon^2 V_1 + \dots$$

and find $V_0 = V_0(S, t')$ only;

then

$$\frac{\partial V_1}{\partial \tau} = \mathcal{L}V_0, \quad \mathcal{L} = \text{zero-div BS operator.}$$

So

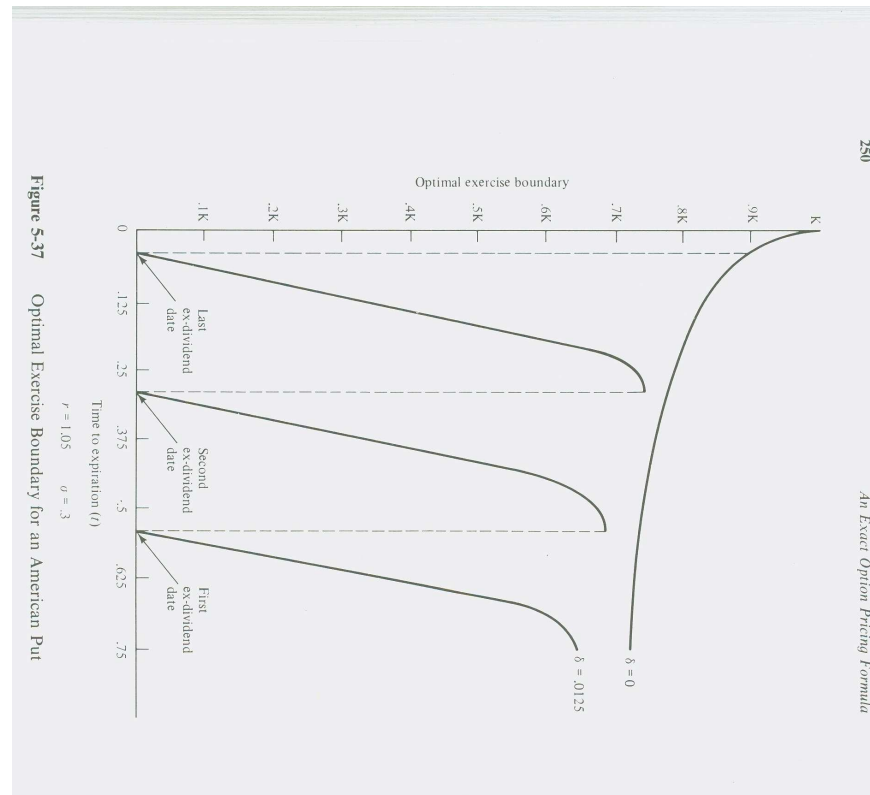
$$V_1 = \tau \mathcal{L}V_0 + F(S, t')$$

and then periodicity plus expanding jump cond'n to $O(\epsilon^2)$ gives

$$\mathcal{L}V_0 = \gamma S \frac{\partial V_0}{\partial S}$$

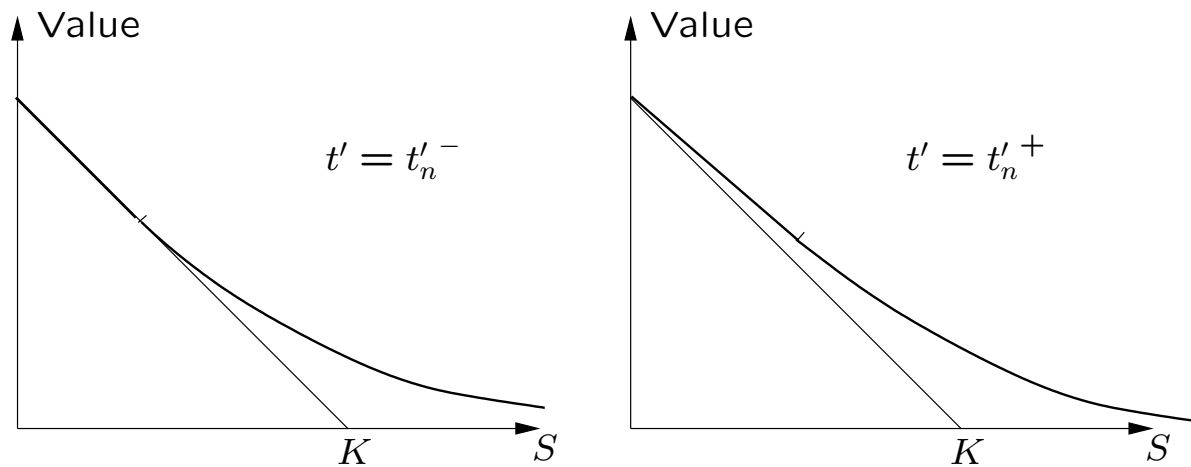
as required.

American option with discrete dividends



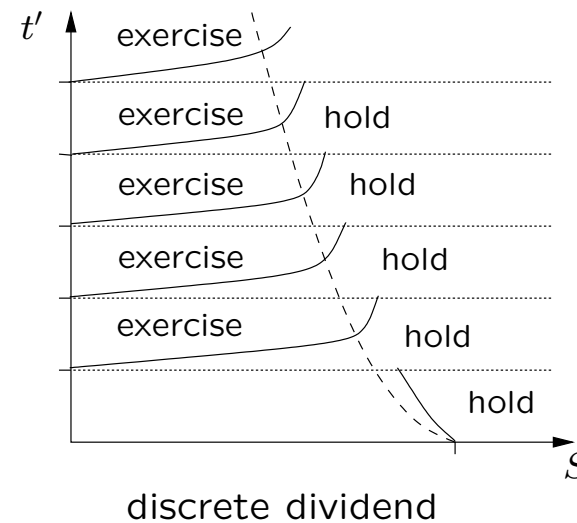
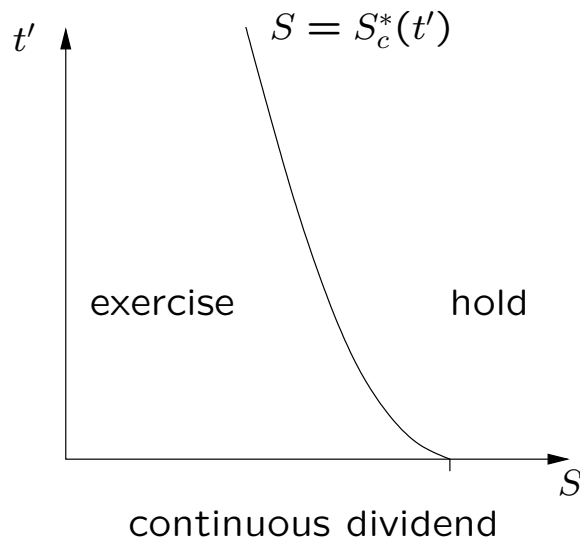
Cox & Rubinstein 1985.

The discrete dividend payment lifts the value function off the payoff:



So the exercise boundary falls to $S = 0$ just after (in backwards time) a dividend date.

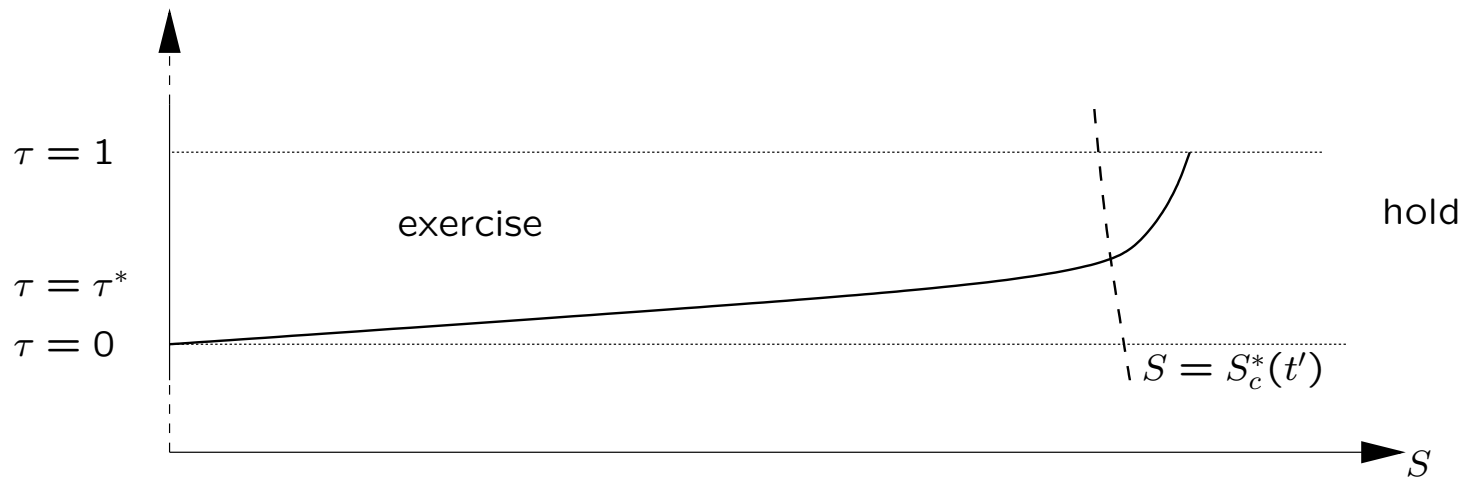
With multiple dividend dates (cf Cox & Rubinstein 1985):



A kind of 'mushy region' homogenised.

Asymptotics of typical inter-dividend period

Put $t' = t'_n + \epsilon^2 \tau$ as before.



Set $V = K - S + W(S, t', \tau)$ and then

$$\frac{\partial W}{\partial t'} + \frac{1}{\epsilon^2} \frac{\partial W}{\partial \tau} = \frac{1}{2} S^2 \frac{\partial^2 W}{\partial S^2} + \rho S \frac{\partial W}{\partial S} - \rho W - \rho K.$$

The free boundary conditions are

$$W = 0, \quad \frac{\partial W}{\partial S} = 0.$$

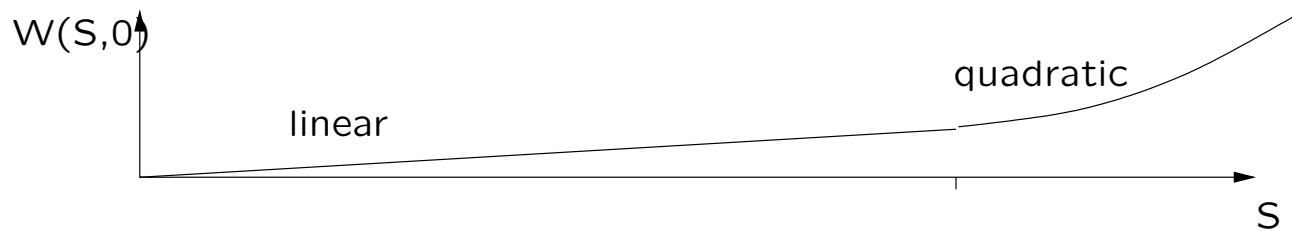
The payoff constraint is

$$W \geq 0.$$

The initial condition (for a periodic solution) is (at leading order)

$$W(S, t', 0) = \begin{cases} \epsilon^2 \gamma S, & 0 < S < S^*(t') \\ \epsilon^2 F(S), & S > S^*(t') \end{cases}$$

where F comes from the outer solution (as above) and it has a (known) second-derivative discontinuity at $S = S^*(t')$. Note it is *linear* for $S < S^*(t')$.



$$\frac{\partial W}{\partial t'} + \frac{1}{\epsilon^2} \frac{\partial W}{\partial \tau} = \frac{1}{2} S^2 \frac{\partial^2 W}{\partial S^2} + \rho S \frac{\partial W}{\partial S} - \rho W - \rho K.$$

The term $-\rho K$ drags W down, but we have $W \geq 0$. Hence a free boundary $S = s^*(\tau)$, at which W vanishes.

Clearly $W \sim \epsilon^2 W_0 + \dots$ and then we have

$$\frac{\partial W_0}{\partial \tau} = -\rho K.$$

Hence

$$W_0 = \gamma S - \rho K \tau$$

for $s^*(\tau) = \rho K \tau / \gamma < S < S^*(t')$, where W_0 vanishes, but only for

$$0 < \tau < \tau^* = \frac{\gamma S^*(t')}{\rho K}.$$

Near $S = s^*(\tau)$ is a small travelling-wave region (as in Stefan) to allow both free boundary conditions to apply.

Meanwhile the curvature jump near $S = S^*(t')$ evolves: put

$$S = S^*(t')(1 + \epsilon x), \quad W = \epsilon^2 w(x, \tau)$$

for an inner region $-\infty < x < \infty$, to get

$$\frac{\partial w}{\partial \tau} = \frac{1}{2} \frac{\partial^2 w}{\partial \tau^2} - \rho K,$$

with initial data having a curvature jump.

Solution is in similarity form.

Transition

This solution only lasts until the free boundary 'sees' the far-field effect of the inner solution. There is a short transition time

$$\tau = \tau^* + O(\epsilon \sqrt{|\log \epsilon|})$$

(determined by 2-term matching...) in which the free boundary behaviour (still the loca-

tion of $W_0 = 0$) goes from

$$s^*(\tau) \sim \rho K \tau / \gamma \quad \text{for } \tau < \tau^* = \gamma S^*(t') / \rho K$$

to

$$s^*(\tau) \sim S^*(t') \left(1 - \epsilon \sqrt{2\tau^*} \left[\sqrt{-\log(\tau - \tau^*)} - \frac{3 \log \sqrt{-\log(\tau - \tau^*)}}{2 \sqrt{-\log(\tau - \tau^*)}} \right] \right)$$

as τ increases away from τ^* .

This is the ‘initial’ condition for a free boundary problem on $x^*(\tau) < x < \infty$, $\tau^* < \tau < 1$. It is the Oxygen Consumption Problem

$$\frac{\partial w}{\partial t} = \frac{1}{2} \frac{\partial^2 w}{\partial x^2} - \rho K, \quad w \geq 0, \quad \tau > \tau^*, \quad x^*(\tau) < x < \infty$$

with initial data

$$w(x, \tau^*) \sim \begin{cases} \text{constant} \times x^{-3} e^{-x^2/2\tau^*} & x \rightarrow -\infty \\ \text{constant} \times x^2 & x \rightarrow +\infty. \end{cases}$$

At $\tau = 1$ it all starts again...

